

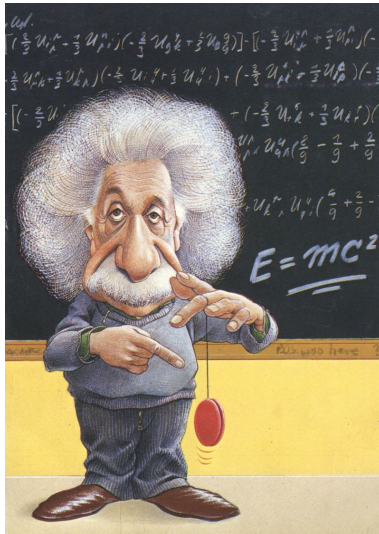
Preventing waves from spreading

Peter Hertel

University of Osnabrück, Germany

Lecture presented at APS, Nankai University, China

March/April 2011



Make it as simple as possible, but not simpler.

Plane waves

- $f(t, \mathbf{x}) \propto e^{i\mathbf{k} \cdot \mathbf{x}} e^{-i\omega t}$

Plane waves

- $f(t, \mathbf{x}) \propto e^{i\mathbf{k} \cdot \mathbf{x}} e^{-i\omega t}$
- wave equation yields $\omega = \omega(\mathbf{k})$

Plane waves

- $f(t, \mathbf{x}) \propto e^{i\mathbf{k} \cdot \mathbf{x}} e^{-i\omega t}$
- wave equation yields $\omega = \omega(\mathbf{k})$
- sound in air: $\omega = v |\mathbf{k}|$

Plane waves

- $f(t, \mathbf{x}) \propto e^{i\mathbf{k} \cdot \mathbf{x}} e^{-i\omega t}$
- wave equation yields $\omega = \omega(\mathbf{k})$
- sound in air: $\omega = v |\mathbf{k}|$
- matter waves (particles): $\omega = \frac{\hbar}{2m} |\mathbf{k}|^2$

Plane waves

- $f(t, \mathbf{x}) \propto e^{i\mathbf{k} \cdot \mathbf{x}} e^{-i\omega t}$
- wave equation yields $\omega = \omega(\mathbf{k})$
- sound in air: $\omega = v |\mathbf{k}|$
- matter waves (particles): $\omega = \frac{\hbar}{2m} |\mathbf{k}|^2$
- light in free space: $\omega = c |\mathbf{k}|$

Wave packets

- Plane wave is an idealization

Wave packets

- Plane wave is an idealization
- Superposition of plane waves, i. e. wave packets

Wave packets

- Plane wave is an idealization
- Superposition of plane waves, i. e. wave packets

- $$f(t, \mathbf{x}) = \int \frac{d^3k}{(2\pi)^3} \phi(\mathbf{k}) e^{i\mathbf{k} \cdot \mathbf{x}} e^{-i\omega(\mathbf{k})t}$$

Wave packets

- Plane wave is an idealization
- Superposition of plane waves, i. e. wave packets

- $$f(t, \mathbf{x}) = \int \frac{d^3k}{(2\pi)^3} \phi(\mathbf{k}) e^{i\mathbf{k} \cdot \mathbf{x}} e^{-i\omega(\mathbf{k})t}$$

- $$\int d^3x |f(t, \mathbf{x})|^2 = \int \frac{d^3k}{(2\pi)^3} |\phi(\mathbf{k})|^2$$

Wave packets

- Plane wave is an idealization
- Superposition of plane waves, i. e. wave packets
- $$f(t, \mathbf{x}) = \int \frac{d^3k}{(2\pi)^3} \phi(\mathbf{k}) e^{i\mathbf{k} \cdot \mathbf{x}} e^{-i\omega(\mathbf{k})t}$$
- $$\int d^3x |f(t, \mathbf{x})|^2 = \int \frac{d^3k}{(2\pi)^3} |\phi(\mathbf{k})|^2$$
- Integral over $|f(t, \mathbf{x})|^2$ does not depend on time

Wave packets

- Plane wave is an idealization
- Superposition of plane waves, i. e. wave packets
- $$f(t, \mathbf{x}) = \int \frac{d^3k}{(2\pi)^3} \phi(\mathbf{k}) e^{i\mathbf{k} \cdot \mathbf{x}} e^{-i\omega(\mathbf{k})t}$$
- $$\int d^3x |f(t, \mathbf{x})|^2 = \int \frac{d^3k}{(2\pi)^3} |\phi(\mathbf{k})|^2$$
- Integral over $|f(t, \mathbf{x})|^2$ does not depend on time
- We normalize it to 1

Location of the wave packet

- $\langle \mathbf{X} \rangle_t = \int d^3x \mathbf{x} |f(t, \mathbf{x})|^2 =$

Location of the wave packet

- $\langle \mathbf{X} \rangle_t = \int d^3x \mathbf{x} |f(t, \mathbf{x})|^2 =$
- $\int \frac{d^3k}{(2\pi)^3} e^{i\omega t} \phi^*(\mathbf{k}) i \nabla_{\mathbf{k}} \phi(\mathbf{k}) e^{-i\omega t} =$

Location of the wave packet

Waves

The Electro-
magnetic
field

Waveguides

Read more

- $\langle \mathbf{X} \rangle_t = \int d^3x \mathbf{x} |f(t, \mathbf{x})|^2 =$
- $\int \frac{d^3k}{(2\pi)^3} e^{i\omega t} \phi^*(\mathbf{k}) i \nabla_{\mathbf{k}} \phi(\mathbf{k}) e^{-i\omega t} =$
- $\int \frac{d^3k}{(2\pi)^3} \phi^*(\mathbf{k}) i \nabla_{\mathbf{k}} \phi(\mathbf{k}) +$

Location of the wave packet

Waves

The Electro-magnetic field

Waveguides

Read more

- $\langle \mathbf{X} \rangle_t = \int d^3x \mathbf{x} |f(t, \mathbf{x})|^2 =$
- $\int \frac{d^3k}{(2\pi)^3} e^{i\omega t} \phi^*(\mathbf{k}) i \nabla_{\mathbf{k}} \phi(\mathbf{k}) e^{-i\omega t} =$
- $\int \frac{d^3k}{(2\pi)^3} \phi^*(\mathbf{k}) i \nabla_{\mathbf{k}} \phi(\mathbf{k}) +$
- $t \int \frac{d^3k}{(2\pi)^3} |\phi(\mathbf{k})|^2 \nabla_{\mathbf{k}} \omega(\mathbf{k})$

Location of the wave packet

Waves

The Electro-magnetic field

Waveguides

Read more

- $\langle \mathbf{X} \rangle_t = \int d^3x \mathbf{x} |f(t, \mathbf{x})|^2 =$
- $\int \frac{d^3k}{(2\pi)^3} e^{i\omega t} \phi^*(\mathbf{k}) i \nabla_{\mathbf{k}} \phi(\mathbf{k}) e^{-i\omega t} =$
- $\int \frac{d^3k}{(2\pi)^3} \phi^*(\mathbf{k}) i \nabla_{\mathbf{k}} \phi(\mathbf{k}) +$
- $t \int \frac{d^3k}{(2\pi)^3} |\phi(\mathbf{k})|^2 \nabla_{\mathbf{k}} \omega(\mathbf{k})$
- $\langle \mathbf{X} \rangle_t = \langle \mathbf{X} \rangle_0 + \mathbf{v} t$

Location of the wave packet

Waves

The Electro-magnetic field

Waveguides

Read more

- $\langle \mathbf{X} \rangle_t = \int d^3x \mathbf{x} |f(t, \mathbf{x})|^2 =$
- $\int \frac{d^3k}{(2\pi)^3} e^{i\omega t} \phi^*(\mathbf{k}) i \nabla_{\mathbf{k}} \phi(\mathbf{k}) e^{-i\omega t} =$
- $\int \frac{d^3k}{(2\pi)^3} \phi^*(\mathbf{k}) i \nabla_{\mathbf{k}} \phi(\mathbf{k}) +$
- $t \int \frac{d^3k}{(2\pi)^3} |\phi(\mathbf{k})|^2 \nabla_{\mathbf{k}} \omega(\mathbf{k})$
- $\langle \mathbf{X} \rangle_t = \langle \mathbf{X} \rangle_0 + \mathbf{v} t$
- $\mathbf{v} = \langle \nabla \omega \rangle = \int \frac{d^3k}{(2\pi)^3} |\phi(\mathbf{k})|^2 \nabla_{\mathbf{k}} \omega(\mathbf{k})$

Location of the wave packet

- $\langle \mathbf{X} \rangle_t = \int d^3x \mathbf{x} |f(t, \mathbf{x})|^2 =$
- $\int \frac{d^3k}{(2\pi)^3} e^{i\omega t} \phi^*(\mathbf{k}) i \nabla_{\mathbf{k}} \phi(\mathbf{k}) e^{-i\omega t} =$
- $\int \frac{d^3k}{(2\pi)^3} \phi^*(\mathbf{k}) i \nabla_{\mathbf{k}} \phi(\mathbf{k}) +$
- $t \int \frac{d^3k}{(2\pi)^3} |\phi(\mathbf{k})|^2 \nabla_{\mathbf{k}} \omega(\mathbf{k})$
- $\langle \mathbf{X} \rangle_t = \langle \mathbf{X} \rangle_0 + \mathbf{v} t$
- $\mathbf{v} = \langle \nabla \omega \rangle = \int \frac{d^3k}{(2\pi)^3} |\phi(\mathbf{k})|^2 \nabla_{\mathbf{k}} \omega(\mathbf{k})$
- group velocity

Spread of the wave packet

- $\langle \mathbf{X}^2 \rangle_t = \int d^3x \mathbf{x}^2 |f(t, \mathbf{x})|^2$

Spread of the wave packet

- $\langle \mathbf{X}^2 \rangle_t = \int d^3x \mathbf{x}^2 |f(t, \mathbf{x})|^2$
- spread $\delta X(t) = \sqrt{\langle \mathbf{X}^2 \rangle_t - \langle \mathbf{X} \rangle_t^2}$

Spread of the wave packet

- $\langle \mathbf{X}^2 \rangle_t = \int d^3x \mathbf{x}^2 |f(t, \mathbf{x})|^2$
- spread $\delta X(t) = \sqrt{\langle \mathbf{X}^2 \rangle_t - \langle \mathbf{X} \rangle_t^2}$
- by a similar calculation as before:

Spread of the wave packet

- $\langle \mathbf{X}^2 \rangle_t = \int d^3x \mathbf{x}^2 |f(t, \mathbf{x})|^2$
- spread $\delta X(t) = \sqrt{\langle \mathbf{X}^2 \rangle_t - \langle \mathbf{X} \rangle_t^2}$
- by a similar calculation as before:
- $\langle \mathbf{X}^2 \rangle_t = \dots + t^2 \langle\langle (\nabla \omega)^2 \rangle\rangle$

Spread of the wave packet

- $\langle \mathbf{X}^2 \rangle_t = \int d^3x \mathbf{x}^2 |f(t, \mathbf{x})|^2$
- spread $\delta X(t) = \sqrt{\langle \mathbf{X}^2 \rangle_t - \langle \mathbf{X} \rangle_t^2}$
- by a similar calculation as before:
- $\langle \mathbf{X}^2 \rangle_t = \dots + t^2 \langle\langle (\nabla \omega)^2 \rangle\rangle$
- for large times t the spread grows as

Spread of the wave packet

- $\langle \mathbf{X}^2 \rangle_t = \int d^3x \mathbf{x}^2 |f(t, \mathbf{x})|^2$
- spread $\delta X(t) = \sqrt{\langle \mathbf{X}^2 \rangle_t - \langle \mathbf{X} \rangle_t^2}$
- by a similar calculation as before:
- $\langle \mathbf{X}^2 \rangle_t = \dots + t^2 \langle\langle (\nabla \omega)^2 \rangle\rangle$
- for large times t the spread grows as
- $\delta X(t) = |t| \sqrt{\langle\langle (\nabla \omega)^2 \rangle\rangle - \langle\langle \nabla \omega \rangle\rangle^2}$

Spread of the wave packet

- $\langle \mathbf{X}^2 \rangle_t = \int d^3x \mathbf{x}^2 |f(t, \mathbf{x})|^2$
- spread $\delta X(t) = \sqrt{\langle \mathbf{X}^2 \rangle_t - \langle \mathbf{X} \rangle_t^2}$
- by a similar calculation as before:
- $\langle \mathbf{X}^2 \rangle_t = \dots + t^2 \langle\langle (\nabla \omega)^2 \rangle\rangle$
- for large times t the spread grows as
- $\delta X(t) = |t| \sqrt{\langle\langle (\nabla \omega)^2 \rangle\rangle - \langle\langle \nabla \omega \rangle\rangle^2}$
- the argument of the square root cannot be negative

Spread of the wave packet

- $\langle \mathbf{X}^2 \rangle_t = \int d^3x \mathbf{x}^2 |f(t, \mathbf{x})|^2$
- spread $\delta X(t) = \sqrt{\langle \mathbf{X}^2 \rangle_t - \langle \mathbf{X} \rangle_t^2}$
- by a similar calculation as before:
- $\langle \mathbf{X}^2 \rangle_t = \dots + t^2 \langle\langle (\nabla \omega)^2 \rangle\rangle$
- for large times t the spread grows as
- $\delta X(t) = |t| \sqrt{\langle\langle (\nabla \omega)^2 \rangle\rangle - \langle\langle \nabla \omega \rangle\rangle^2}$
- the argument of the square root cannot be negative
- Wave packets finally spread out. . .

Spread of the wave packet

- $\langle \mathbf{X}^2 \rangle_t = \int d^3x \mathbf{x}^2 |f(t, \mathbf{x})|^2$
- spread $\delta X(t) = \sqrt{\langle \mathbf{X}^2 \rangle_t - \langle \mathbf{X} \rangle_t^2}$
- by a similar calculation as before:
- $\langle \mathbf{X}^2 \rangle_t = \dots + t^2 \langle\langle (\nabla \omega)^2 \rangle\rangle$
- for large times t the spread grows as
- $\delta X(t) = |t| \sqrt{\langle\langle (\nabla \omega)^2 \rangle\rangle - \langle\langle \nabla \omega \rangle\rangle^2}$
- the argument of the square root cannot be negative
- Wave packets finally spread out. . .
- . . . if the medium is homogeneous .

Electromagnetic field

- The electromagnetic field is defined by its action on charged particles

Electromagnetic field

- The electromagnetic field is defined by its action on charged particles
- $\dot{\mathbf{p}} = q\{\mathbf{E} + \mathbf{v} \times \mathbf{B}\}$

Electromagnetic field

- The electromagnetic field is defined by its action on charged particles
- $\dot{\mathbf{p}} = q\{\mathbf{E} + \mathbf{v} \times \mathbf{B}\}$
- location \mathbf{x} , velocity \mathbf{v} , momentum \mathbf{p}

Electromagnetic field

- The electromagnetic field is defined by its action on charged particles
- $\dot{\mathbf{p}} = q\{\mathbf{E} + \mathbf{v} \times \mathbf{B}\}$
- location \mathbf{x} , velocity \mathbf{v} , momentum \mathbf{p}
- charge q , electric field strength \mathbf{E} , magnetic induction \mathbf{B}

Electromagnetic field

- The electromagnetic field is defined by its action on charged particles
- $\dot{\mathbf{p}} = q\{\mathbf{E} + \mathbf{v} \times \mathbf{B}\}$
- location \mathbf{x} , velocity \mathbf{v} , momentum \mathbf{p}
- charge q , electric field strength \mathbf{E} , magnetic induction \mathbf{B}
- The electromagnetic field is generated by a distribution of charged particles

Electromagnetic field

- The electromagnetic field is defined by its action on charged particles
- $\dot{\mathbf{p}} = q\{\mathbf{E} + \mathbf{v} \times \mathbf{B}\}$
- location \mathbf{x} , velocity \mathbf{v} , momentum \mathbf{p}
- charge q , electric field strength \mathbf{E} , magnetic induction \mathbf{B}
- The electromagnetic field is generated by a distribution of charged particles
- charge density ρ , current density \mathbf{j}

Electromagnetic field

- The electromagnetic field is defined by its action on charged particles
- $\dot{\mathbf{p}} = q\{\mathbf{E} + \mathbf{v} \times \mathbf{B}\}$
- location \mathbf{x} , velocity \mathbf{v} , momentum \mathbf{p}
- charge q , electric field strength \mathbf{E} , magnetic induction \mathbf{B}
- The electromagnetic field is generated by a distribution of charged particles
- charge density ρ , current density \mathbf{j}
- Maxwell's equations

Electromagnetic field

- The electromagnetic field is defined by its action on charged particles
- $\dot{\mathbf{p}} = q\{\mathbf{E} + \mathbf{v} \times \mathbf{B}\}$
- location \mathbf{x} , velocity \mathbf{v} , momentum \mathbf{p}
- charge q , electric field strength \mathbf{E} , magnetic induction \mathbf{B}
- The electromagnetic field is generated by a distribution of charged particles
- charge density ρ , current density \mathbf{j}
- Maxwell's equations
- $\mathbf{div} \mathbf{D} = \rho, \mathbf{div} \mathbf{B} = 0$

Electromagnetic field

- The electromagnetic field is defined by its action on charged particles
- $\dot{\mathbf{p}} = q\{\mathbf{E} + \mathbf{v} \times \mathbf{B}\}$
- location \mathbf{x} , velocity \mathbf{v} , momentum \mathbf{p}
- charge q , electric field strength \mathbf{E} , magnetic induction \mathbf{B}
- The electromagnetic field is generated by a distribution of charged particles
- charge density ρ , current density \mathbf{j}
- Maxwell's equations
- $\operatorname{div} \mathbf{D} = \rho, \operatorname{div} \mathbf{B} = 0$
- $\operatorname{curl} \mathbf{H} = \mathbf{j} + \dot{\mathbf{D}}, \operatorname{curl} \mathbf{E} = -\dot{\mathbf{B}}$

Electromagnetic field

- The electromagnetic field is defined by its action on charged particles
- $\dot{\mathbf{p}} = q\{\mathbf{E} + \mathbf{v} \times \mathbf{B}\}$
- location \mathbf{x} , velocity \mathbf{v} , momentum \mathbf{p}
- charge q , electric field strength \mathbf{E} , magnetic induction \mathbf{B}
- The electromagnetic field is generated by a distribution of charged particles
- charge density ρ , current density \mathbf{j}
- Maxwell's equations
- $\operatorname{div} \mathbf{D} = \rho, \operatorname{div} \mathbf{B} = 0$
- $\operatorname{curl} \mathbf{H} = \mathbf{j} + \dot{\mathbf{D}}, \operatorname{curl} \mathbf{E} = -\dot{\mathbf{B}}$
- linear Medium: $\mathbf{D} = \epsilon\epsilon_0\mathbf{E}, \mathbf{B} = \mu\mu_0\mathbf{H}$

Preventing waves from spreading

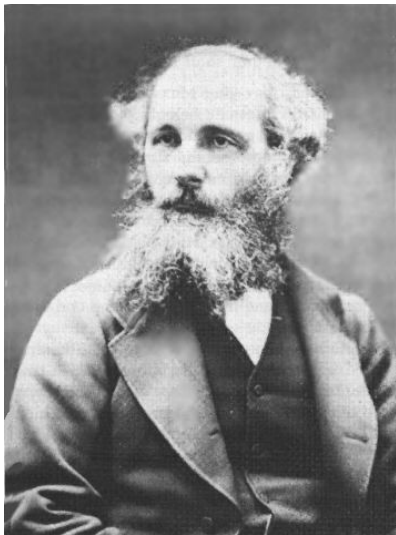
Peter Hertel

Waves

The Electro-magnetic field

Waveguides

Read more



James Clerk Maxwell, 1831-1879

- no charges, no currents: $\rho = 0, \mathbf{j} = 0$

- no charges, no currents: $\rho = 0, \mathbf{j} = 0$
- no magnetic properties: $\mu = 1$

- no charges, no currents: $\rho = 0, \mathbf{j} = 0$
- no magnetic properties: $\mu = 1$
- need to study fields $\propto e^{-i\omega t}$ only

- no charges, no currents: $\rho = 0, \mathbf{j} = 0$
- no magnetic properties: $\mu = 1$
- need to study fields $\propto e^{-i\omega t}$ only
- $\nabla \cdot \epsilon \mathbf{E} = 0, \text{div } \mathbf{H} = 0$

- no charges, no currents: $\rho = 0, \mathbf{j} = 0$
- no magnetic properties: $\mu = 1$
- need to study fields $\propto e^{-i\omega t}$ only
- $\nabla \cdot \mathbf{E} = 0, \operatorname{div} \mathbf{H} = 0$
- $\operatorname{curl} \mathbf{H} = -i\omega\epsilon_0 \epsilon \mathbf{E}, \operatorname{curl} \mathbf{E} = i\omega\mu_0 \mathbf{H}$

- no charges, no currents: $\rho = 0, \mathbf{j} = 0$
- no magnetic properties: $\mu = 1$
- need to study fields $\propto e^{-i\omega t}$ only
- $\nabla \cdot \mathbf{E} = 0, \operatorname{div} \mathbf{H} = 0$
- $\operatorname{curl} \mathbf{H} = -i\omega\epsilon_0 \mathbf{E}, \operatorname{curl} \mathbf{E} = i\omega\mu_0 \mathbf{H}$
- With $\epsilon_0\mu_0c^2 = 1$ and $k_0 = \omega/c$:

- no charges, no currents: $\rho = 0, \mathbf{j} = 0$
- no magnetic properties: $\mu = 1$
- need to study fields $\propto e^{-i\omega t}$ only
- $\nabla \cdot \epsilon \mathbf{E} = 0, \operatorname{div} \mathbf{H} = 0$
- $\operatorname{curl} \mathbf{H} = -i\omega\epsilon_0 \epsilon \mathbf{E}, \operatorname{curl} \mathbf{E} = i\omega\mu_0 \mathbf{H}$
- With $\epsilon_0\mu_0 c^2 = 1$ and $k_0 = \omega/c$:
- **$\operatorname{curl} \operatorname{curl} \mathbf{E} = k_0^2 \epsilon \mathbf{E}$**

- no charges, no currents: $\rho = 0, \mathbf{j} = 0$
- no magnetic properties: $\mu = 1$
- need to study fields $\propto e^{-i\omega t}$ only
- $\nabla \cdot \epsilon \mathbf{E} = 0, \operatorname{div} \mathbf{H} = 0$
- $\operatorname{curl} \mathbf{H} = -i\omega\epsilon_0 \epsilon \mathbf{E}, \operatorname{curl} \mathbf{E} = i\omega\mu_0 \mathbf{H}$
- With $\epsilon_0\mu_0 c^2 = 1$ and $k_0 = \omega/c$:
- **$\operatorname{curl} \operatorname{curl} \mathbf{E} = k_0^2 \epsilon \mathbf{E}$**
- equivalent

- no charges, no currents: $\rho = 0, \mathbf{j} = 0$
- no magnetic properties: $\mu = 1$
- need to study fields $\propto e^{-i\omega t}$ only
- $\nabla \cdot \mathbf{E} = 0, \operatorname{div} \mathbf{H} = 0$
- $\operatorname{curl} \mathbf{H} = -i\omega\epsilon_0 \epsilon \mathbf{E}, \operatorname{curl} \mathbf{E} = i\omega\mu_0 \mathbf{H}$
- With $\epsilon_0\mu_0 c^2 = 1$ and $k_0 = \omega/c$:
- **$\operatorname{curl} \operatorname{curl} \mathbf{E} = k_0^2 \epsilon \mathbf{E}$**
- equivalent
- **$\operatorname{curl} \epsilon^{-1} \operatorname{curl} \mathbf{H} = k_0^2 \mathbf{H}$**

- no charges, no currents: $\rho = 0, \mathbf{j} = 0$
- no magnetic properties: $\mu = 1$
- need to study fields $\propto e^{-i\omega t}$ only
- $\nabla \cdot \epsilon \mathbf{E} = 0, \operatorname{div} \mathbf{H} = 0$
- $\operatorname{curl} \mathbf{H} = -i\omega\epsilon_0 \epsilon \mathbf{E}, \operatorname{curl} \mathbf{E} = i\omega\mu_0 \mathbf{H}$
- With $\epsilon_0\mu_0 c^2 = 1$ and $k_0 = \omega/c$:
- **$\operatorname{curl} \operatorname{curl} \mathbf{E} = k_0^2 \epsilon \mathbf{E}$**
- equivalent
- **$\operatorname{curl} \epsilon^{-1} \operatorname{curl} \mathbf{H} = k_0^2 \mathbf{H}$**
- $\epsilon \mathbf{E}$ and \mathbf{H} are automatically divergence free

Waveguides

- Spreading of light is unavoidable if the medium is homogeneous

Waveguides

- Spreading of light is unavoidable if the medium is homogeneous
- Therefore, the medium must be inhomogeneous if light is to be guided

Waveguides

- Spreading of light is unavoidable if the medium is homogeneous
- Therefore, the medium must be inhomogeneous if light is to be guided
- permittivity profile $\epsilon = \epsilon(\mathbf{x})$

Waveguides

- Spreading of light is unavoidable if the medium is homogeneous
- Therefore, the medium must be inhomogeneous if light is to be guided
- permittivity profile $\epsilon = \epsilon(\mathbf{x})$
- Non-constant imaginary part: microwave guides, coaxial cables

Waveguides

- Spreading of light is unavoidable if the medium is homogeneous
- Therefore, the medium must be inhomogeneous if light is to be guided
- permittivity profile $\epsilon = \epsilon(\mathbf{x})$
- Non-constant imaginary part: microwave guides, coaxial cables
- ϵ real and non-constant: dielectric waveguides

Read more

- Lecture notes are deposited at `ftp://202.113.31.42`

Read more

- Lecture notes are deposited at `ftp://202.113.31.42`
- Change directory to `/temp/peter.hertel/2011-03`

Read more

- Lecture notes are deposited at `ftp://202.113.31.42`
- Change directory to `/temp/peter.hertel/2011-03`
- `dwg.pdf` - Dielectric Waveguides

Read more

- Lecture notes are deposited at `ftp://202.113.31.42`
- Change directory to `/temp/peter.hertel/2011-03`
- `dwg.pdf` - Dielectric Waveguides
- `basics.pdf` - this lecture

Read more

- Lecture notes are deposited at `ftp://202.113.31.42`
- Change directory to `/temp/peter.hertel/2011-03`
- `dwg.pdf` - Dielectric Waveguides
- `basics.pdf` - this lecture
- `planar.pdf` - next lecture