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Permittivity tensor

Maxwell's equations

Isotropic medium

Birefringence

Absorption

Drude model

Crystal optics

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• Crystal optics investigates the propagation of plane waves in a homogeneous medium

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- Crystal optics investigates the propagation of plane waves in a homogeneous medium
- The susceptibility tensor is real and symmetric

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- three eigenvalues equal: isotropic medium

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- three eigenvalues equal: isotropic medium
- only two are equal: uniaxial medium

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• for a sufficiently weak light wave, the polarization is linear in the electric field strength

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Linear medium

- for a sufficiently weak light wave, the polarization is linear in the electric field strength
- however retarded, but local

$$P_i(t, \boldsymbol{x}) = \epsilon_0 \int_0^\infty \mathrm{d}\tau \, G_{ij}(\tau) \, E_j(t - \tau, \boldsymbol{x})$$

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• Einstein's summation convention: sum over j from 1 to 3

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$$F(t) = \int \frac{\mathrm{d}\omega}{2\pi} f(\omega) \, e^{-\mathrm{i}\omega t}$$

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- the Fourier transform of the dispacement is

 $d_i(\omega, \boldsymbol{x}) = \epsilon_0 \, \epsilon_{ij}(\omega) \, e_j(\omega, \boldsymbol{x})$

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- the Fourier transform of the dispacement is

$$d_i(\omega, \boldsymbol{x}) = \epsilon_0 \, \epsilon_{ij}(\omega) \, e_j(\omega, \boldsymbol{x})$$

• $\epsilon_{ij}(\omega) = \delta_{ij} + \chi_{ij}(\omega)$ where $\chi_{ij} = \tilde{G}_{ij}$

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Refraction and absorption

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• decompose

$$\epsilon_{ij} = \epsilon'_{ij} + \mathrm{i}\,\epsilon''_{ij}$$

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Refraction and absorption

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• decompose

$$\epsilon_{ij} = \epsilon'_{ij} + i \, \epsilon''_{ij}$$

$$\epsilon_{ij}' = \frac{\epsilon_{ij} + \epsilon_{ji}^*}{2}$$

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Refraction and absorption

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• decompose

$$\epsilon_{ij} = \epsilon'_{ij} + \mathrm{i}\,\epsilon''_{ij}$$

refractive part

$$\epsilon_{ij}' = \frac{\epsilon_{ij} + \epsilon_{ji}^*}{2}$$

$$\epsilon_{ij}'' = \frac{\epsilon_{ij} - \epsilon_{ji}}{2i}$$

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• both are Hermitian: $A_{ij} = A_{ji}^*$

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- we shall see later while absorptive part causes absorption

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Refraction and absorption

• decompose

$$\epsilon_{ij} = \epsilon'_{ij} + i \, \epsilon''_{ij}$$

refractive part

$$\epsilon_{ij}' = \frac{\epsilon_{ij} + \epsilon_{ji}^*}{2}$$

• absorptive part

$$\epsilon_{ij}'' = \frac{\epsilon_{ij} - \epsilon_{ji}^*}{2\mathrm{i}}$$

- both are Hermitian: $A_{ij} = A_{ji}^*$
- we shall see later while absorptive part causes absorption
- i. e. the conversion of field energy into internal energy

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• assume $\epsilon_{ij}^{\prime\prime}(\omega)\approx 0$ for the frequencies under discussion

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• assume $\epsilon_{ij}''(\omega)\approx 0$ for the frequencies under discussion • then $\epsilon_{ij}=\epsilon_{ij}'$ is hermitian

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Transparent medium

- assume $\epsilon_{ij}^{\prime\prime}(\omega)\approx 0$ for the frequencies under discussion
- then $\epsilon_{ij} = \epsilon'_{ij}$ is hermitian
- recall Onsager's relations

$$\epsilon_{ij}(\omega; \boldsymbol{\mathcal{E}}, \boldsymbol{\mathcal{B}}) = \epsilon_{ji}(\omega; \boldsymbol{\mathcal{E}}, -\boldsymbol{\mathcal{B}})$$

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• where $\boldsymbol{\mathcal{E}}, \boldsymbol{\mathcal{B}}$ are static external fields

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- without external induction field

$$\epsilon_{ji} = \epsilon_{ij} = \epsilon_{ij}^*$$

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- without external induction field

 $\epsilon_{ji} = \epsilon_{ij} = \epsilon_{ij}^*$

• no absorption, no external induction: ϵ_{ij} is a real symmetric matrix

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- without external induction field

 $\epsilon_{ji} = \epsilon_{ij} = \epsilon_{ij}^*$

- no absorption, no external induction: ϵ_{ij} is a real symmetric matrix
- ϵ_{ij} can be diagonalized by an orthogonal matrix

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there is a Cartesian coordinate system such that

$$\epsilon_{ij} = \left(\begin{array}{ccc} \epsilon^1 & 0 & 0\\ 0 & \epsilon^2 & 0\\ 0 & 0 & \epsilon^3 \end{array}\right)$$

Main axes

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$$\epsilon_{ij} = \left(\begin{array}{ccc} \epsilon^1 & 0 & 0 \\ 0 & \epsilon^2 & 0 \\ 0 & 0 & \epsilon^3 \end{array} \right)$$

There are three cases **1** $\epsilon^1 = \epsilon^2 = \epsilon^3$ isotropic medium (glass)

Main axes

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There are three cases (1) $\epsilon^1 = \epsilon^2 = \epsilon^3$ isotropic medium (glass) (2) $\epsilon^1 = \epsilon^2 \neq \epsilon^3$ uniaxial (LiNbO₃)

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there is a Cartesian coordinate system such that

$$\epsilon_{ij} = \left(\begin{array}{ccc} \epsilon^1 & 0 & 0 \\ 0 & \epsilon^2 & 0 \\ 0 & 0 & \epsilon^3 \end{array} \right)$$

There are three cases

 $\begin{array}{l} \bullet \ \epsilon^1 = \epsilon^2 = \epsilon^3 \\ \bullet \ \epsilon^1 = \epsilon^2 \neq \epsilon^3 \\ \bullet \ \epsilon^1 < \epsilon^2 < \epsilon^3 \\ \bullet^1 < \epsilon^2 \\ \bullet^1 < \epsilon^2$

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• assume plane waves with fixed angular frequency

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- assume plane waves with fixed angular frequency
- all fields are of the form

$$F(t, \boldsymbol{x}) = f e^{-i\omega t} e^{i\boldsymbol{k}\cdot\boldsymbol{x}}$$

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$$\boldsymbol{k} \times \boldsymbol{e} = \omega \mu_0 \boldsymbol{h}$$

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$$\boldsymbol{k} \times \boldsymbol{h} = -\omega \epsilon_0 \epsilon \boldsymbol{e}$$

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•
$$\boldsymbol{k} imes \boldsymbol{e} = \omega \mu_0 \boldsymbol{h}$$

•
$$\boldsymbol{k} imes \boldsymbol{h} = -\omega \epsilon_0 \, \epsilon \, \boldsymbol{e}$$

• With
$$c = 1/\sqrt{\epsilon_0 \mu_0}$$

$$(\boldsymbol{k} imes \boldsymbol{k} imes \boldsymbol{e})_i = -rac{\omega^2}{c^2} \epsilon_{ij} e_j$$

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$$(m{k} imes m{k} imes m{e})_i = -rac{\omega^2}{c^2} \epsilon_{ij} e_j$$

• $k_0 = \omega/c$ vacuum wave number

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$$(m{k} imes m{k} imes m{e})_i = -rac{\omega^2}{c^2} \epsilon_{ij} e_j$$

- $k_0 = \omega/c$ vacuum wave number
- $m{k} = n k_0 \hat{m{k}}$ refractive index n, propagation direction $\hat{m{k}}$

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$$(m{k} imes m{k} imes m{e})_i = -rac{\omega^2}{c^2} \epsilon_{ij} e_j$$

- $k_0 = \omega/c$ vacuum wave number
- $\boldsymbol{k} = nk_0\hat{\boldsymbol{k}}$ refractive index n, propagation direction $\hat{\boldsymbol{k}}$
- $e = e\hat{e}$ polarization vector \hat{e}

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$$(m{k} imes m{k} imes m{e})_i = -rac{\omega^2}{c^2} \epsilon_{ij} e_j$$

- $k_0 = \omega/c$ vacuum wave number
- $\boldsymbol{k} = nk_0\hat{\boldsymbol{k}}$ refractive index n, propagation direction $\hat{\boldsymbol{k}}$
- $e = e\hat{e}$ polarization vector \hat{e}
- to be solved is the mode equation $n^{2}(\hat{\boldsymbol{k}} \times \hat{\boldsymbol{k}} \times \hat{\boldsymbol{e}})_{i} = -\epsilon_{ij}\hat{\boldsymbol{e}}_{j}$

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•
$$\boldsymbol{c} \times \boldsymbol{b} \times \boldsymbol{a} = (\boldsymbol{c} \cdot \boldsymbol{a})\boldsymbol{b} - (\boldsymbol{c} \cdot \boldsymbol{b})\boldsymbol{a}$$

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$$\boldsymbol{c} \times \boldsymbol{b} \times \boldsymbol{a} = (\boldsymbol{c} \cdot \boldsymbol{a})\boldsymbol{b} - (\boldsymbol{c} \cdot \boldsymbol{b})\boldsymbol{a}$$

• $\hat{\boldsymbol{k}} \times \hat{\boldsymbol{k}} \times \hat{\boldsymbol{e}} = (\hat{\boldsymbol{k}} \cdot \hat{\boldsymbol{e}})\hat{\boldsymbol{k}} - \hat{\boldsymbol{e}}$

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$$\boldsymbol{c} \times \boldsymbol{b} \times \boldsymbol{a} = (\boldsymbol{c} \cdot \boldsymbol{a})\boldsymbol{b} - (\boldsymbol{c} \cdot \boldsymbol{b})\boldsymbol{a}$$

• $\hat{\boldsymbol{k}} \times \hat{\boldsymbol{k}} \times \hat{\boldsymbol{e}} = (\hat{\boldsymbol{k}} \cdot \hat{\boldsymbol{e}})\hat{\boldsymbol{k}} - \hat{\boldsymbol{e}}$

• mode equation can be written as $n^2(\hat{e} - (\hat{k} \cdot \hat{e})\hat{k})_i = \epsilon_{ij}\hat{e}_j$

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$$\boldsymbol{c} \times \boldsymbol{b} \times \boldsymbol{a} = (\boldsymbol{c} \cdot \boldsymbol{a})\boldsymbol{b} - (\boldsymbol{c} \cdot \boldsymbol{b})\boldsymbol{a}$$

• $\hat{\boldsymbol{k}} \times \hat{\boldsymbol{k}} \times \hat{\boldsymbol{e}} = (\hat{\boldsymbol{k}} \cdot \hat{\boldsymbol{e}})\hat{\boldsymbol{k}} - \hat{\boldsymbol{e}}$

• mode equation can be written as

$$n^2(\hat{\boldsymbol{e}} - (\hat{\boldsymbol{k}} \cdot \hat{\boldsymbol{e}})\hat{\boldsymbol{k}})_i = \epsilon_{ij}\hat{\boldsymbol{e}}_j$$

• no solution for
$$\hat{m{k}} \parallel \hat{m{e}}$$

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- Birefringence
- Absorption
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Maxwell's equations ctd.

•
$$\boldsymbol{c} \times \boldsymbol{b} \times \boldsymbol{a} = (\boldsymbol{c} \cdot \boldsymbol{a})\boldsymbol{b} - (\boldsymbol{c} \cdot \boldsymbol{b})\boldsymbol{a}$$

• $\hat{\boldsymbol{k}} \times \hat{\boldsymbol{k}} \times \hat{\boldsymbol{e}} = (\hat{\boldsymbol{k}} \cdot \hat{\boldsymbol{e}})\hat{\boldsymbol{k}} - \hat{\boldsymbol{e}}$

• mode equation can be written as
$$n^2(\hat{m{e}}-(\hat{m{k}}\cdot\hat{m{e}})\hat{m{k}})_i=\epsilon_{ij}\hat{m{e}}_j$$

- no solution for $\hat{m{k}} \parallel \hat{m{e}}$
- therefore $\hat{m{k}} \perp \hat{m{e}}$

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Maxwell's equations ctd.

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$$\mathbf{c} \times \mathbf{b} \times \mathbf{a} = (\mathbf{c} \cdot \mathbf{a})\mathbf{b} - (\mathbf{c} \cdot \mathbf{b})\mathbf{a}$$

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$$\hat{\boldsymbol{k}} \times \hat{\boldsymbol{k}} \times \hat{\boldsymbol{e}} = (\hat{\boldsymbol{k}} \cdot \hat{\boldsymbol{e}})\hat{\boldsymbol{k}} - \hat{\boldsymbol{e}}$$

- mode equation can be written as $n^2(\hat{e} - (\hat{k} \cdot \hat{e})\hat{k})_i = \epsilon_{ij}\hat{e}_j$
- no solution for $\hat{m{k}} \parallel \hat{m{e}}$
- therefore $rac{\hat{k} \perp \hat{e}}{k}$
- electromagnetic plane waves in a homogeneous medium are alway transversally polarized

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Optically isotropic medium

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Optically isotropic medium

• if
$$\epsilon^1 = \epsilon^2 = \epsilon^3 = \epsilon$$

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Optically isotropic medium

• if
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• we find
$$\epsilon_{ij} = \epsilon \, \delta_{ij}$$

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Optically isotropic medium

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• this is true for an arbitrary Cartesian coordinate system

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Optically isotropic medium

- if $\epsilon^1 = \epsilon^2 = \epsilon^3 = \epsilon$
- we find $\epsilon_{ij} = \epsilon \, \delta_{ij}$
- this is true for an arbitrary Cartesian coordinate system
- we say the medium is optically isotropic

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Optically isotropic medium

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$$n^2(\hat{\boldsymbol{e}} - (\hat{\boldsymbol{k}} \cdot \hat{\boldsymbol{e}})\,\hat{\boldsymbol{k}}) = \epsilon \hat{\boldsymbol{e}}$$

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$$\hat{k} \perp \hat{e}$$
: $n = \sqrt{\epsilon}$

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- $\hat{m{h}} = \hat{m{k}} imes \hat{m{e}}$

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- any polarization \hat{e} is allowed.
- any orthogonal propagation direction $\hat{m{k}}$ is allowed
- $\hat{m{h}} = \hat{m{k}} imes \hat{m{e}}$
- \hat{k} , \hat{e} , \hat{h} is right handed set of orthogonal unit vectors

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Optical axis

- Two eigenvalues of ϵ_{ij} are equal, we call them ordinary

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Optical axis

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• Two eigenvalues of ϵ_{ij} are equal, we call them ordinary

• the extraordinary eigenvalue is different

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Optical axis

- Two eigenvalues of ϵ_{ij} are equal, we call them ordinary
- the extraordinary eigenvalue is different
- it belongs to the optical axis (here \hat{z})

$$\epsilon_{ij} = \left(\begin{array}{ccc} \epsilon^{\mathrm{o}} & 0 & 0 \\ 0 & \epsilon^{\mathrm{o}} & 0 \\ 0 & 0 & \epsilon^{\mathrm{e}} \end{array}\right)$$

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• ordinary beam is polarized \perp optical axis

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• ordinary beam is polarized \perp optical axis

•
$$\hat{\boldsymbol{e}} = \cos\phi\,\hat{\boldsymbol{x}} + \sin\phi\,\hat{\boldsymbol{y}}$$
, $\hat{\boldsymbol{k}} = \hat{\boldsymbol{z}}$, $n^{\mathrm{o}} = \sqrt{\epsilon^{\mathrm{o}}}$

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- extraordinary beam is polarized || optical axis

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- extraordinary beam is polarized || optical axis

•
$$\hat{m{e}}=\hat{m{z}}$$
, $\hat{m{k}}=\coslpha\,\hat{m{x}}+\sinlpha\,\hat{m{y}}$, $n^{
m e}=\sqrt{\epsilon^{
m e}}$

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• what happens if a beam is polarized neither parallel nor perpendicular to optical axis?

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- what happens if a beam is polarized neither parallel nor perpendicular to optical axis?
- e. g. if it is unpolarized

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- what happens if a beam is polarized neither parallel nor perpendicular to optical axis?
- e. g. if it is unpolarized
- when entering the medium it splits into an ordinary and an extraordinary beam

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- what happens if a beam is polarized neither parallel nor perpendicular to optical axis?
- e. g. if it is unpolarized
- when entering the medium it splits into an ordinary and an extraordinary beam
- which propagate with different refractive index

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- what happens if a beam is polarized neither parallel nor perpendicular to optical axis?
- e. g. if it is unpolarized
- when entering the medium it splits into an ordinary and an extraordinary beam
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- they will leave the the medium at different locations

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- when entering the medium it splits into an ordinary and an extraordinary beam
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- double refraction, or birefringence

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Birefringence

- what happens if a beam is polarized neither parallel nor perpendicular to optical axis?
- e. g. if it is unpolarized
- when entering the medium it splits into an ordinary and an extraordinary beam
- which propagate with different refractive index
- they will leave the the medium at different locations
- being polarized
- double refraction, or birefringence
- calcite (at which birefringence was discovered)

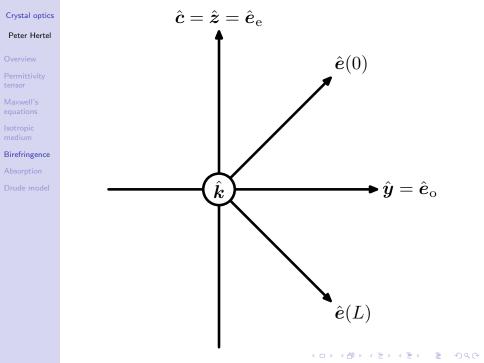
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- what happens if a beam is polarized neither parallel nor perpendicular to optical axis?
- e. g. if it is unpolarized
- when entering the medium it splits into an ordinary and an extraordinary beam
- which propagate with different refractive index
- they will leave the the medium at different locations
- being polarized
- double refraction, or birefringence
- calcite (at which birefringence was discovered)
- elastooptics



Crystal optics	A 45 degree polarized wave enters the crystal and leaves it at
Peter Hertel	-45 degrees polarization.
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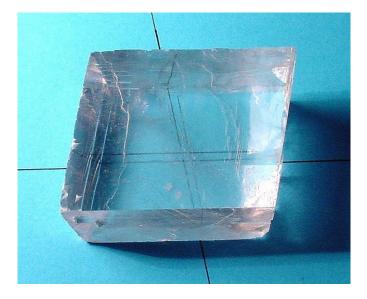
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Birefringence, or double refraction, by calcite

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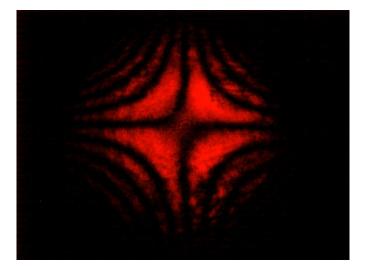
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Normally isotropic polymers become birefringent when stressed. Observed with a polarizer.

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• Only the ordinary beam may be unpolarized

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- Only the ordinary beam may be unpolarized
- optically biaxial media have three different eigenvalues of permittivity tensor ϵ_{ij}

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• Only the ordinary beam may be unpolarized

- optically biaxial media have three different eigenvalues of permittivity tensor ϵ_{ij}
- correspondingly three orthogonal directions

$$\epsilon_{ij} = \left(\begin{array}{ccc} \epsilon^1 & 0 & 0 \\ 0 & \epsilon^2 & 0 \\ 0 & 0 & \epsilon^3 \end{array} \right)$$

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- beams polarized along these axes propagate with different refractive indexes $n^a=\sqrt{\epsilon^a}$

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- beams polarized along these axes propagate with different refractive indexes $n^a=\sqrt{\epsilon^a}$
- rather difficult to show that there are two optical axes

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• for simplicity, assume isotropic medium



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• for simplicity, assume isotropic medium

• $\epsilon = \epsilon' + i\epsilon''$

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• for simplicity, assume isotropic medium

- $\epsilon = \epsilon' + i\epsilon''$
- nearly transparent medium, $\epsilon^{\prime\prime} \ll \epsilon^\prime$

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• for simplicity, assume isotropic medium

- $\epsilon = \epsilon' + i\epsilon''$
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• recall
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- $\hat{m{k}} \perp \hat{m{e}}$ remains true

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$$\bar{n} = \sqrt{\epsilon' + i\epsilon''}$$
 is complex

• With
$$n = \sqrt{\epsilon'}$$
 one may write $\bar{n} \approx n + \mathrm{i} \frac{\epsilon''}{2n}$

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- $\epsilon = \epsilon' + i\epsilon''$
- nearly transparent medium, $\epsilon'' \ll \epsilon'$

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• with $lpha=\epsilon^{\prime\prime}k_0/n$ and $z=\hat{m k}\cdotm x$ one finds

$$\boldsymbol{E}(t,\boldsymbol{x}) = \boldsymbol{E}(0,0) e^{-\mathrm{i}\omega t} e^{\mathrm{i}nk_0 z} e^{-\alpha z/2}$$

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• with $\alpha = \epsilon'' k_0/n$ and $z = \hat{k} \cdot x$ one finds

$$\boldsymbol{E}(t,\boldsymbol{x}) = \boldsymbol{E}(0,0) e^{-\mathrm{i}\omega t} e^{\mathrm{i}nk_0 z} e^{-\alpha z/2}$$

• $S \propto |\mathbf{E}|^2$, $S(z) = S(0) e^{-\alpha z}$

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$$S \propto |\mathbf{E}|^2$$
, $S(z) = S(0) e^{-\alpha z}$

• α is absorption constant

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Overview

Permittivity tensor

Maxwell's equations

lsotropic medium

Birefringence

Absorption

Drude model

Drude model

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• consider typical electron with mass m and charge q=-e

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• consider typical electron with mass m and charge q = -e

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- $oldsymbol{x}$ is deviation from equilibrium position
- damped harmonic oscillatation

$$m(\ddot{\boldsymbol{x}} + \Gamma \dot{\boldsymbol{x}} + \Omega^2 \boldsymbol{x}) = q\boldsymbol{E}$$

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$$\chi(\omega) = \chi(0) \frac{\Omega^2}{\Omega^2 - \omega^2 - i\Gamma\omega}$$

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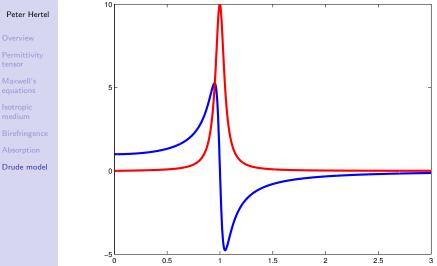
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static susceptibility

$$\chi(0) = \frac{Nq^2}{m\Omega^2\epsilon_0}$$





Real (blue) and imaginary part (red) of susceptibility $\chi(\omega)$ relative to $\chi(0)$ over ω/Ω . $\Gamma/\Omega = 0.1$