Properties of susceptibility

Peter Hertel

University of Osnabrück, Germany

April 1, 2010

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

$$\blacktriangleright P_i(t,\mathbf{x}) = \sum_j \int d\tau \int d^3 \boldsymbol{\xi} \ \theta(\tau) \Gamma_{ij} u, \boldsymbol{\xi}) E_j(t-\tau,\mathbf{x}-\boldsymbol{\xi})$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 - のへぐ

$$P_i(t, \mathbf{x}) = \sum_j \int d\tau \int d^3 \boldsymbol{\xi} \ \theta(\tau) \Gamma_{ij} u, \boldsymbol{\xi} E_j(t - \tau, \mathbf{x} - \boldsymbol{\xi})$$

$$\hat{P}_i(\omega, \mathbf{q}) = \epsilon_0 \sum_j \chi_{ij}(\omega, \mathbf{q}) \hat{E}_j(\omega, \mathbf{q})$$

$$P_i(t, \mathbf{x}) = \sum_j \int d\tau \int d^3 \boldsymbol{\xi} \ \theta(\tau) \Gamma_{ij} u, \boldsymbol{\xi}) E_j(t - \tau, \mathbf{x} - \boldsymbol{\xi})$$

$$\hat{P}_i(\omega, \mathbf{q}) = \epsilon_0 \sum_j \chi_{ij}(\omega, \mathbf{q}) \hat{E}_j(\omega, \mathbf{q})$$

The susceptibility is given by

$$\chi_{ij}(\omega,\mathbf{q}) = \frac{1}{\epsilon_0} \int d\tau \, e^{\,i\omega\tau} \,\,\theta(\tau) \int d^3\xi \, e^{\,-i\mathbf{q}\cdot\boldsymbol{\xi}} \,\Gamma_{ij}(\tau,\boldsymbol{\xi})$$

$$P_i(t, \mathbf{x}) = \sum_j \int d\tau \int d^3 \boldsymbol{\xi} \ \theta(\tau) \Gamma_{ij} u, \boldsymbol{\xi} E_j(t - \tau, \mathbf{x} - \boldsymbol{\xi})$$

$$\hat{P}_i(\omega, \mathbf{q}) = \epsilon_0 \sum_j \chi_{ij}(\omega, \mathbf{q}) \hat{E}_j(\omega, \mathbf{q})$$

The susceptibility is given by

$$\chi_{ij}(\omega,\mathbf{q}) = \frac{1}{\epsilon_0} \int d\tau \, e^{i\omega\tau} \, \theta(\tau) \int d^3\xi \, e^{-i\mathbf{q}\cdot\boldsymbol{\xi}} \, \Gamma_{ij}(\tau,\boldsymbol{\xi})$$

Let's concentrate on time

$$P_i(t, \mathbf{x}) = \sum_j \int d\tau \int d^3 \boldsymbol{\xi} \ \theta(\tau) \Gamma_{ij} u, \boldsymbol{\xi}) E_j(t - \tau, \mathbf{x} - \boldsymbol{\xi})$$

$$\hat{P}_i(\omega, \mathbf{q}) = \epsilon_0 \sum_j \chi_{ij}(\omega, \mathbf{q}) \hat{E}_j(\omega, \mathbf{q})$$

The susceptibility is given by

$$\chi_{ij}(\omega,\mathbf{q}) = \frac{1}{\epsilon_0} \int d\tau \, e^{\,i\omega\tau} \,\,\theta(\tau) \int d^3\xi \, e^{\,-i\mathbf{q}\cdot\boldsymbol{\xi}} \,\Gamma_{ij}(\tau,\boldsymbol{\xi})$$

Let's concentrate on time

Introduce

$$\Gamma_{ij}(\tau,\mathbf{q}) = \int d^3\xi \, e^{-i\mathbf{q}\cdot\boldsymbol{\xi}} \, \Gamma_{ij}(\tau,\boldsymbol{\xi})$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

$$P_i(t, \mathbf{x}) = \sum_j \int d\tau \int d^3 \boldsymbol{\xi} \ \theta(\tau) \Gamma_{ij} u, \boldsymbol{\xi}) E_j(t - \tau, \mathbf{x} - \boldsymbol{\xi})$$

$$\hat{P}_i(\omega, \mathbf{q}) = \epsilon_0 \sum_j \chi_{ij}(\omega, \mathbf{q}) \hat{E}_j(\omega, \mathbf{q})$$

The susceptibility is given by

$$\chi_{ij}(\omega,\mathbf{q}) = \frac{1}{\epsilon_0} \int d\tau \, e^{\,i\omega\tau} \,\,\theta(\tau) \int d^3\xi \, e^{\,-i\mathbf{q}\cdot\boldsymbol{\xi}} \,\Gamma_{ij}(\tau,\boldsymbol{\xi})$$

Let's concentrate on time

Introduce

$$\Gamma_{ij}(\tau,\mathbf{q}) = \int d^3 \boldsymbol{\xi} \, e^{-i\mathbf{q} \cdot \boldsymbol{\xi}} \, \Gamma_{ij}(\tau,\boldsymbol{\xi})$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

• Drop indexes i, j and the wave-vector **q** and forget about ϵ_0

$$P_i(t, \mathbf{x}) = \sum_j \int d\tau \int d^3 \boldsymbol{\xi} \ \theta(\tau) \Gamma_{ij} u, \boldsymbol{\xi}) E_j(t - \tau, \mathbf{x} - \boldsymbol{\xi})$$

$$\hat{P}_i(\omega, \mathbf{q}) = \epsilon_0 \sum_j \chi_{ij}(\omega, \mathbf{q}) \hat{E}_j(\omega, \mathbf{q})$$

The susceptibility is given by

$$\chi_{ij}(\omega,\mathbf{q}) = \frac{1}{\epsilon_0} \int d\tau \, e^{\,i\omega\tau} \,\,\theta(\tau) \int d^3\xi \, e^{\,-i\mathbf{q}\cdot\boldsymbol{\xi}} \,\Gamma_{ij}(\tau,\boldsymbol{\xi})$$

Let's concentrate on time

Introduce

$$\Gamma_{ij}(\tau,\mathbf{q}) = \int d^3\boldsymbol{\xi} \, e^{-i\mathbf{q}\cdot\boldsymbol{\xi}} \, \Gamma_{ij}(\tau,\boldsymbol{\xi})$$

Drop indexes i, j and the wave-vector q and forget about e₀

We shall study next

$$\chi(\omega) = \int d\tau \, e^{\,i\omega\tau} \, \theta(\tau) \, \Gamma(\tau)$$

▲□▶ ▲圖▶ ★ 国▶ ★ 国▶ - 国 - のへで

• $\chi = \chi(\omega)$ is the Fourier transform of a causal function

• $\chi = \chi(\omega)$ is the Fourier transform of a causal function

• namely of $f(t)=\theta(t)\Gamma(t)$

• $\chi = \chi(\omega)$ is the Fourier transform of a causal function

- namely of $f(t)=\theta(t)\Gamma(t)$
- which therefore may be written as $f(t) = \theta(t)f(t)$

• $\chi = \chi(\omega)$ is the Fourier transform of a causal function

- namely of $f(t) = \theta(t)\Gamma(t)$
- which therefore may be written as $f(t) = \theta(t)f(t)$
- hence, according to the convolution theorem:

$$\hat{f}(\omega) = \int \frac{du}{2\pi} \hat{ heta}(\omega - u) \hat{f}(u)$$

• $\chi = \chi(\omega)$ is the Fourier transform of a causal function

- namely of $f(t) = \theta(t)\Gamma(t)$
- which therefore may be written as $f(t) = \theta(t)f(t)$
- hence, according to the convolution theorem:

$$\hat{f}(\omega) = \int \frac{du}{2\pi} \hat{ heta}(\omega - u) \hat{f}(u)$$

• With $\epsilon > 0$, $\epsilon \rightarrow 0$ the Fourier transform of θ is

$$\hat{\theta}(\omega) = \frac{1}{\epsilon - i\omega}$$

• $\chi = \chi(\omega)$ is the Fourier transform of a causal function

- namely of $f(t) = \theta(t)\Gamma(t)$
- which therefore may be written as $f(t) = \theta(t)f(t)$
- hence, according to the convolution theorem:

$$\hat{f}(\omega) = \int \frac{du}{2\pi} \hat{ heta}(\omega - u) \hat{f}(u)$$

• With $\epsilon > 0$, $\epsilon \rightarrow 0$ the Fourier transform of θ is

$$\hat{ heta}(\omega) = rac{1}{\epsilon - i\omega}$$

Hence

$$\chi(\omega) = \int \frac{du}{2\pi} \frac{\chi(u)}{\epsilon - i(\omega - u)} = \frac{1}{2\pi i} \int \frac{du \, \chi(u)}{u - \omega + i\epsilon}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ ◆□ ◆ ◇◇◇

▶ The subsceptibility is related with itself:

$$\chi(\omega) = \int \frac{du}{2\pi} \frac{\chi(u)}{\epsilon - i(\omega - u)} = \frac{1}{2\pi i} \int \frac{du \ \chi(u)}{u - \omega - i\epsilon}$$

◆□▶ ◆□▶ ◆∃▶ ◆∃▶ = のへで

▶ The subsceptibility is related with itself:

$$\chi(\omega) = \int \frac{du}{2\pi} \frac{\chi(u)}{\epsilon - i(\omega - u)} = \frac{1}{2\pi i} \int \frac{du \ \chi(u)}{u - \omega - i\epsilon}$$

• Split into real and imaginary part: $\chi = \chi' + i\chi''$

▶ The subsceptibility is related with itself:

$$\chi(\omega) = \int \frac{du}{2\pi} \frac{\chi(u)}{\epsilon - i(\omega - u)} = \frac{1}{2\pi i} \int \frac{du \ \chi(u)}{u - \omega - i\epsilon}$$

• Split into real and imaginary part: $\chi = \chi' + i\chi''$

 \blacktriangleright recall that χ is the Fourier transform of a real function

▶ The subsceptibility is related with itself:

$$\chi(\omega) = \int \frac{du}{2\pi} \frac{\chi(u)}{\epsilon - i(\omega - u)} = \frac{1}{2\pi i} \int \frac{du \ \chi(u)}{u - \omega - i\epsilon}$$

- Split into real and imaginary part: $\chi = \chi' + i\chi''$
- \blacktriangleright recall that χ is the Fourier transform of a real function

•
$$\chi(\omega) = \chi^*(-\omega)$$

▶ The subsceptibility is related with itself:

$$\chi(\omega) = \int \frac{du}{2\pi} \frac{\chi(u)}{\epsilon - i(\omega - u)} = \frac{1}{2\pi i} \int \frac{du \ \chi(u)}{u - \omega - i\epsilon}$$

- Split into real and imaginary part: $\chi = \chi' + i\chi''$
- \blacktriangleright recall that χ is the Fourier transform of a real function

•
$$\chi(\omega) = \chi^*(-\omega)$$

This boils down to

$$\chi'(\omega) = \frac{2}{\pi} \int_0^\infty du \, \frac{u \, \chi''(u)}{u^2 - \omega^2}$$

The subsceptibility is related with itself:

$$\chi(\omega) = \int \frac{du}{2\pi} \frac{\chi(u)}{\epsilon - i(\omega - u)} = \frac{1}{2\pi i} \int \frac{du \ \chi(u)}{u - \omega - i\epsilon}$$

- Split into real and imaginary part: $\chi = \chi' + i\chi''$
- \blacktriangleright recall that χ is the Fourier transform of a real function
- $\chi(\omega) = \chi^*(-\omega)$
- This boils down to

$$\chi'(\omega) = \frac{2}{\pi} \int_0^\infty du \, \frac{u \, \chi''(u)}{u^2 - \omega^2}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

The subsceptibility is related with itself:

$$\chi(\omega) = \int \frac{du}{2\pi} \frac{\chi(u)}{\epsilon - i(\omega - u)} = \frac{1}{2\pi i} \int \frac{du \ \chi(u)}{u - \omega - i\epsilon}$$

- Split into real and imaginary part: $\chi = \chi' + i\chi''$
- \blacktriangleright recall that χ is the Fourier transform of a real function
- $\chi(\omega) = \chi^*(-\omega)$
- This boils down to

$$\chi'(\omega) = \frac{2}{\pi} \int_0^\infty du \, \frac{u \, \chi''(u)}{u^2 - \omega^2}$$

- principal value integral understood
- note that only positive (angular) frequencies are involved

The subsceptibility is related with itself:

$$\chi(\omega) = \int \frac{du}{2\pi} \frac{\chi(u)}{\epsilon - i(\omega - u)} = \frac{1}{2\pi i} \int \frac{du \ \chi(u)}{u - \omega - i\epsilon}$$

- Split into real and imaginary part: $\chi = \chi' + i\chi''$
- \blacktriangleright recall that χ is the Fourier transform of a real function

•
$$\chi(\omega) = \chi^*(-\omega)$$

This boils down to

$$\chi'(\omega) = \frac{2}{\pi} \int_0^\infty du \, \frac{u \, \chi''(u)}{u^2 - \omega^2}$$

- principal value integral understood
- note that only positive (angular) frequencies are involved
- χ'' describes absorption. Why must it be positive?

•
$$A(t) = U_{-t}AU_t$$
 is a process



•
$$A(t) = U_{-t}AU_t$$
 is a process

• note that < A(t) > does not vary with time

•
$$A(t) = U_{-t}AU_t$$
 is a process

- note that < A(t) > does not vary with time
- Define the correlation function for two processes

$$K(A, B; \tau) = rac{}{2} -$$

(ロ)、(型)、(E)、(E)、 E) の(の)

•
$$A(t) = U_{-t}AU_t$$
 is a process

- note that <A(t)> does not vary with time
- Define the correlation function for two processes

$${\cal K}(A,B; au)=rac{<\!A(t+ au)B(t)+B(t)A(t+ au)>}{2}-<\!A\!><\!B\!>$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへで

note that time t does not appear on the left hand side

•
$$A(t) = U_{-t}AU_t$$
 is a process

- note that < A(t) > does not vary with time
- Define the correlation function for two processes

$${\cal K}(A,B; au)=rac{<\!A(t+ au)B(t)+B(t)A(t+ au)>}{2}-<\!A\!><\!B\!>$$

- note that time t does not appear on the left hand side
- ▶ note that the correlation function K = K(A, B; τ) is real if A and B are observables

•
$$A(t) = U_{-t}AU_t$$
 is a process

- note that <A(t)> does not vary with time
- Define the correlation function for two processes

$${\cal K}(A,B; au)=rac{<\!A(t+ au)B(t)+B(t)A(t+ au)>}{2}-<\!A\!><\!B\!>$$

- note that time t does not appear on the left hand side
- ▶ note that the correlation function K = K(A, B; τ) is real if A and B are observables

• $K(A, A; \tau)$ is called an auto-correlation function

•
$$A(t) = U_{-t}AU_t$$
 is a process

- note that <A(t)> does not vary with time
- Define the correlation function for two processes

$${\cal K}(A,B; au)=rac{<\!A(t+ au)B(t)+B(t)A(t+ au)>}{2}-<\!A\!><\!B\!>$$

- note that time t does not appear on the left hand side
- ▶ note that the correlation function K = K(A, B; τ) is real if A and B are observables

- $K(A, A; \tau)$ is called an auto-correlation function
- ► K(A, B, 0) describes the correlation between A and B

Let us insert Fourier-transforms

$$A(t+\tau) = + \int \frac{d\omega}{2\pi} e^{-i\omega\(t+\tau\)} \hat{A}\(\omega\)$$
$$B(t) = + \int \frac{d\omega'}{2\pi} e^{i\omega't} \hat{B}^{\dagger}(\omega')$$

Let us insert Fourier-transforms

$$A(t+\tau) = + \int \frac{d\omega}{2\pi} e^{-i\omega\(t+\tau\)} \hat{A}\(\omega\)$$
$$B(t) = + \int \frac{d\omega'}{2\pi} e^{i\omega't} \hat{B}^{\dagger}(\omega')$$

There will be a dependency on t unless

$$rac{<\hat{\mathcal{A}}(\omega)\hat{B}^{\dagger}(\omega^{\,\prime})+\hat{B}^{\dagger}(\omega^{\,\prime})\hat{\mathcal{A}}(\omega)>}{2}=2\pi\delta(\omega-\omega^{\,\prime})\,\mathcal{S}(\omega)$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

Let us insert Fourier-transforms

$$A(t+\tau) = + \int \frac{d\omega}{2\pi} e^{-i\omega\(t+\tau\)} \hat{A}\(\omega\)$$
$$B(t) = + \int \frac{d\omega'}{2\pi} e^{i\omega't} \hat{B}^{\dagger}(\omega')$$

There will be a dependency on t unless

$$rac{<\hat{\mathcal{A}}(\omega)\hat{B}^{\dagger}(\omega^{\,\prime})+\hat{B}^{\dagger}(\omega^{\,\prime})\hat{\mathcal{A}}(\omega)>}{2}=2\pi\delta(\omega-\omega^{\,\prime})\,\mathcal{S}(\omega)$$

With this spectral density we may write

$$K(AB; \tau) = \int \frac{d\omega}{2\pi} e^{-i\omega\tau} S(A, B; \omega)$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

Let us insert Fourier-transforms

$$egin{aligned} A(t+ au) = & + \int rac{d\omega}{2\pi} \, e^{-i\omega\(t+ au\)} \, \hat{A}\(\omega\) \ B\(t\) = & + \int rac{d\omega'}{2\pi} \, e^{i\omega't} \, \hat{B}^{\dagger}\(\omega'\) \end{aligned}$$

There will be a dependency on t unless

$$rac{<\hat{\mathcal{A}}(\omega)\hat{B}^{\dagger}(\omega^{\,\prime})+\hat{B}^{\dagger}(\omega^{\,\prime})\hat{\mathcal{A}}(\omega)>}{2}=2\pi\delta(\omega-\omega^{\,\prime})\,\mathcal{S}(\omega)$$

With this spectral density we may write

$$K(AB; \tau) = \int \frac{d\omega}{2\pi} e^{-i\omega\tau} S(A, B; \omega)$$

The Wiener Khinchin theorem says

$$K(AA; \tau) = \int \frac{d\omega}{2\pi} e^{-i\omega\tau} S(A, A; \omega)$$

with $S(AA; \omega) \geq 0$



Norbert Wiener

- 4 同 ト 4 三 ト 4 三



Aleksandr Yakovlevich Khinchin

Define

$$A(z) = e^{-\frac{i}{\hbar}zH}Ae^{\frac{i}{\hbar}zH}$$

Define

$$A(z) = e^{-\frac{i}{\hbar}zH}Ae^{\frac{i}{\hbar}zH}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

Both the Gibbs state as well as the waiting operator are exponentials of the energy H

Define

$$A(z) = e^{-\frac{i}{\hbar}zH}Ae^{\frac{i}{\hbar}zH}$$

- Both the Gibbs state as well as the waiting operator are exponentials of the energy H
- We work out

$$A(z) e^{-\beta H} = e^{-\beta H} e^{\beta H} A(z) e^{-\beta H}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

Define

$$A(z) = e^{-\frac{i}{\hbar}zH}Ae^{\frac{i}{\hbar}zH}$$

- Both the Gibbs state as well as the waiting operator are exponentials of the energy H
- We work out

$$A(z) e^{-\beta H} = e^{-\beta H} e^{\beta H} A(z) e^{-\beta H}$$

• i.e. $A(z)G = GA(z - i\hbar\beta)$

Define

$$A(z) = e^{-\frac{i}{\hbar}zH}Ae^{\frac{i}{\hbar}zH}$$

- Both the Gibbs state as well as the waiting operator are exponentials of the energy H
- We work out

$$A(z) e^{-\beta H} = e^{-\beta H} e^{\beta H} A(z) e^{-\beta H}$$

• i.e.
$$A(z)G = GA(z - i\hbar\beta)$$

Multiply with B from the right and apply the trace

Define

$$A(z) = e^{-\frac{i}{\hbar}zH}Ae^{\frac{i}{\hbar}zH}$$

- Both the Gibbs state as well as the waiting operator are exponentials of the energy H
- We work out

$$A(z) e^{-\beta H} = e^{-\beta H} e^{\beta H} A(z) e^{-\beta H}$$

• i.e.
$$A(z)G = GA(z - i\hbar\beta)$$

- Multiply with B from the right and apply the trace
- KMS formula $\langle BA(z) \rangle = \langle A(z i\hbar\beta)B \rangle$

Define

$$A(z) = e^{-\frac{i}{\hbar}zH}Ae^{\frac{i}{\hbar}zH}$$

- Both the Gibbs state as well as the waiting operator are exponentials of the energy H
- We work out

$$A(z) e^{-\beta H} = e^{-\beta H} e^{\beta H} A(z) e^{-\beta H}$$

• i.e.
$$A(z)G = GA(z - i\hbar\beta)$$

- Multiply with B from the right and apply the trace
- KMS formula $\langle BA(z) \rangle = \langle A(z i\hbar\beta)B \rangle$
- where $\beta = 1/kBT$



Ryogo Kubo



Paul Martin

▲□▶ ▲□▶ ▲目▶ ▲目▶

æ



Julian Schwinger

◆□▶ ◆舂▶ ◆臣▶ ◆臣▶

æ

Compare

$$\Gamma(AB; \tau) = < \frac{i}{\hbar}[A(\tau), B] >$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

• Compare
$$\Gamma(AB;\tau) = <\frac{i}{\hbar}[A(\tau),B]>$$

with

$$\mathcal{K}(AB;\tau) = \frac{\langle A(\tau)B + BA(\tau) \rangle}{2} - \langle A \rangle \langle B \rangle$$

Compare
$$\Gamma(AB; au)=<rac{i}{\hbar}[A(au),B]>$$

with

$$K(AB;\tau) = \frac{\langle A(\tau)B + BA(\tau) \rangle}{2} - \langle A \rangle \langle B \rangle$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

• Define $f(\tau) = \langle A(\tau)B \rangle - \langle A \rangle \langle B \rangle$

• Compare
$$\Gamma(AB;\tau) = <\frac{i}{\hbar}[A(\tau),B]>$$

with

$$K(AB;\tau) = \frac{\langle A(\tau)B + BA(\tau) \rangle}{2} - \langle A \rangle \langle B \rangle$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

• Define
$$f(\tau) = \langle A(\tau)B \rangle - \langle A \rangle \langle B \rangle$$

 This function can be continued uniquely into the complex plane

• Compare
$$\Gamma(AB; \tau) = <rac{i}{\hbar}[A(\tau), B]>$$

with

$$K(AB;\tau) = \frac{\langle A(\tau)B + BA(\tau) \rangle}{2} - \langle A \rangle \langle B \rangle$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

• Define
$$f(\tau) = \langle A(\tau)B \rangle - \langle A \rangle \langle B \rangle$$

- This function can be continued uniquely into the complex plane
- Applying the KMS formula yields

• Compare
$$\Gamma(AB; \tau) = < \frac{i}{\hbar} [A(\tau), B] >$$

with

$$K(AB;\tau) = \frac{\langle A(\tau)B + BA(\tau) \rangle}{2} - \langle A \rangle \langle B \rangle$$

• Define
$$f(\tau) = \langle A(\tau)B \rangle - \langle A \rangle \langle B \rangle$$

- This function can be continued uniquely into the complex plane
- Applying the KMS formula yields
- $\blacktriangleright < BA(\tau) > < A > < B > = f(\tau i\hbar\beta)$

Response function can be written as

$$\Gamma(AB;\tau) = \frac{i}{\hbar} \{f(\tau) - f(\tau - i\hbar\beta)\}$$

Response function can be written as

$$\Gamma(AB;\tau) = \frac{i}{\hbar} \{f(\tau) - f(\tau - i\hbar\beta)\}$$

Time correlation function is

$$K(AB;\tau) = \frac{1}{2} \{ f(\tau) + f(\tau - i\hbar\beta) \}$$

Response function can be written as

$$\Gamma(AB; \tau) = rac{i}{\hbar} \{f(\tau) - f(\tau - i\hbar\beta)\}$$

Time correlation function is

$$K(AB; \tau) = \frac{1}{2} \{ f(\tau) + f(\tau - i\hbar\beta) \}$$

Generalized susceptibility

$$\gamma(AB;\omega) = \int_0^\infty d\tau \ e^{i\omega\tau} \,\Gamma(AB;\tau)$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

Response function can be written as

$$\Gamma(AB; \tau) = rac{i}{\hbar} \{f(\tau) - f(\tau - i\hbar\beta)\}$$

Time correlation function is

$$K(AB; \tau) = \frac{1}{2} \{ f(\tau) + f(\tau - i\hbar\beta) \}$$

Generalized susceptibility

$$\gamma(AB;\omega) = \int_0^\infty d au \ e^{i\omega au} \, \Gamma(AB; au)$$

Fluctuation-Dissipation theorem of Callen and Welton

$$\frac{\gamma(AB;\omega)-\gamma(BA;\omega)^*}{2i}=S(A,B;\omega)\frac{1}{\hbar}\tanh\frac{\beta\hbar\omega}{2}$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

Response function can be written as

$$\Gamma(AB; \tau) = rac{i}{\hbar} \{ f(\tau) - f(\tau - i\hbar\beta) \}$$

Time correlation function is

$$K(AB; \tau) = \frac{1}{2} \{ f(\tau) + f(\tau - i\hbar\beta) \}$$

Generalized susceptibility

$$\gamma(AB;\omega) = \int_0^\infty d\tau \ e^{i\omega\tau} \, \Gamma(AB;\tau)$$

Fluctuation-Dissipation theorem of Callen and Welton

$$\frac{\gamma(AB;\omega)-\gamma(BA;\omega)^*}{2i}=S(A,B;\omega)\frac{1}{\hbar}\tanh\frac{\beta\hbar\omega}{2}$$



J. Callen

< □ > < □ > < □ > < □ > < □ > < □ > = □